

1 **Suppression of cold air outbreaks over the interior of North America in a**  
2 **warmer climate**

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7 ABSTRACT: In spite of the mean warming trend over the last few decades and its amplification in  
8 the Arctic, some studies have found no robust decline or even a slight increase in wintertime cold  
9 air outbreaks over North America. But fossil evidence from warmer paleoclimate periods indicates  
10 that the interior of North America never dropped below freezing even in the depths of winter,  
11 which implies that the maintenance of cold air outbreaks is unlikely to continue indefinitely with  
12 future warming. To identify key mechanisms affecting cold air outbreaks and understand how and  
13 why they will change in a warmer climate, we examine the development of North American cold  
14 air outbreaks in both a pre-industrial and a roughly  $8\times\text{CO}_2$  scenario using the Community Earth  
15 System Model, CESM2. We observe a sharp drop-off in the wintertime temperature distribution  
16 at the freezing temperature, suppressing below-freezing conditions in the warmer climate and  
17 above-freezing conditions in the pre-industrial case. The disappearance of Arctic sea ice and  
18 loss of the near-surface temperature inversion dramatically decrease the availability of below-  
19 freezing air in source regions. Using an air parcel trajectory analysis, we demonstrate a remarkable  
20 similarity in both the dynamics and diabatic effects acting on cold air masses in the two climate  
21 scenarios. Diabatic temperature evolution along cold air outbreak trajectories is a competition  
22 between cooling from longwave radiation and warming from boundary layer mixing. Surprisingly,  
23 while both diabatic effects strengthen in the warmer climate, the balance remains the same, with a  
24 net cooling of about  $-6$  K over 10 days.

25 SIGNIFICANCE STATEMENT: We compare a pre-industrial climate scenario to a much warmer  
26 climate circa the year 2300 under high emissions to understand the physical processes that influence  
27 the coldest wintertime temperatures and how they will change with warming. We find that enhanced  
28 warming in the Arctic, and particularly over the Arctic Ocean due to the loss of wintertime sea  
29 ice, dramatically reduces the availability of cold air to be swept into North America. By tracing  
30 these cold air masses as they travel, we also find that they experience the same total amount of  
31 cooling in the much warmer climate as they did in the pre-industrial climate, even though many of  
32 the individual heating and cooling processes have gotten stronger.

### 33 **1. Introduction**

34 There is a pressing need to understand trends in modern cold air outbreaks. Despite the overall  
35 warming trend in global mean temperature over the past few decades (Arias et al. 2021), there  
36 is some disagreement over whether there has been a corresponding warming in North American  
37 cold air outbreaks. Some studies have found a decline in wintertime cold air outbreak frequency,  
38 intensity, or duration (Hankes and Walsh 2011; Robeson et al. 2014; Screen 2014; Grotjahn et al.  
39 2016; Oldenborgh et al. 2019; Smith and Sheridan 2020), while others have found no robust trend  
40 (Walsh et al. 2001; Portis et al. 2006; Westby et al. 2013) or even a slight increase in some regions  
41 (Liu et al. 2012; Cohen et al. 2014). If there are mechanisms acting to maintain or even strengthen  
42 cold air mass formation in spite of the warming trend, we would do well to identify them so that  
43 we can better predict where and when, or even if, they will disappear as warming continues.

44 Model projections of changes in North American cold air outbreaks over the next few decades  
45 are also subject to significant variability, further solidifying the need for a better mechanistic  
46 understanding. A study by Vavrus et al. (2006) on the change in cold air outbreaks by the middle  
47 of the twenty-first century identified large spatial and inter-model variability, with some regions  
48 experiencing a total loss of cold air outbreaks and others no significant change. Identifying the  
49 factors responsible for variability would improve confidence in predictions, particularly over the  
50 next few decades where models disagree on the trend. A series of studies have implicated a  
51 sensitivity to remote sea ice and snow cover conditions (Cohen and Entekhabi 1999; Gong et al.  
52 2003; Vavrus 2007; Cohen et al. 2014; Vavrus et al. 2017), which can vary dramatically from  
53 year to year. Preconditioning of the surface that cold air masses pass over can also influence air

54 mass evolution, as a recent history of cold conditions or snow cover can enhance the ability of the  
55 land surface to act as a heat sink in the development of the coldest air masses (Ellis and Leathers  
56 1998; Gao et al. 2015; Hartig et al. 2023). Inter-model and spatial variability could also result  
57 from competing factors, such as the cooling from longwave radiation balanced by warming from  
58 turbulent convection identified in (Hartig et al. 2023), in a competition that has not fully tipped to  
59 either side at the observed level of warming.

60 The question of cold air outbreaks in a warmer climate is also motivated by the Eocene, a warm  
61 climate period that persisted from 56 to 34 million years ago. The Eocene presents a challenge to  
62 our understanding of cold air outbreaks in a warmer climate due to strong evidence for the complete  
63 suppression of below-freezing temperatures over North America. Estimated CO<sub>2</sub> levels during the  
64 Eocene range from 600 to over 1,500 ppm, depending on the time period and the proxy used  
65 (Beerling and Royer 2011), and the global mean surface temperature was 10 to 16 K higher than  
66 in the pre-industrial climate (Inglis et al. 2020). But wintertime continental interiors were much,  
67 much warmer than they are today. Fossils of frost-intolerant species from the Eocene, including  
68 palms, turtles, and crocodiles, have been found in the central Great Plains and Rocky Mountains  
69 of North America (Hutchison 1982; Wing and Greenwood 1993; Markwick 1994; Greenwood and  
70 Wing 1995). The presence of these species implies no more than a day at a time below freezing  
71 and an absolute minimum temperature above  $-10^{\circ}\text{C}$  (Wing and Greenwood 1993; Greenwood and  
72 Wing 1995; Hyland et al. 2018) with a cold month mean temperature of at least  $4^{\circ}\text{C}$  (Markwick  
73 1994, 1998) and as high as  $13^{\circ}\text{C}$  (Hutchison 1982). Today, those regions have a cold month mean  
74 of  $-4^{\circ}\text{C}$  and experience over 100 days per year below freezing with typical wintertime minimums  
75 of  $-30^{\circ}\text{C}$  to  $-40^{\circ}\text{C}$  (NWS 2023), indicating that cold extremes in the Eocene were suppressed by  
76 a factor of 2–3 relative to the warming of the wintertime mean.

77 While multiple lines of fossil evidence point to above-freezing temperatures over Eocene con-  
78 tinental interiors, climate models have consistently struggled to simulate conditions that match  
79 the proxies. A recent Eocene Model Intercomparison Project noted difficulties matching even the  
80 global mean surface temperature over many models in the project (Lunt et al. 2012). Models  
81 that can match high- and mid-latitude temperatures produce tropics that are too warm (Shellito  
82 et al. 2003), or conversely match the tropics but not higher latitudes (Heinemann et al. 2009). An  
83 older CCSM3 model produced reasonable global temperatures but required 4560 ppm CO<sub>2</sub> to do

84 so, which the authors interpret as unrealistically low climate sensitivity rather than a challenge to  
85 CO<sub>2</sub> proxies (Huber and Caballero 2011; Caballero and Huber 2013). The previous generation of  
86 CESM produced a good match for mean annual temperatures across latitudes in the Eocene (Zhu  
87 et al. 2019) but did not look at seasonal means or minimums. The struggle across many models to  
88 simulate a feasible Eocene climate suggests that a deeper understanding of the physical processes  
89 that enhance or suppress cold extremes will help to identify what is going wrong when models  
90 attempt to produce above-freezing continental temperatures in an Eocene-like climate.

91 In this study, we compare a pre-industrial climate scenario to a much warmer climate to un-  
92 derstand the physical processes that influence the coldest wintertime temperatures. The warmer  
93 climate corresponds to an extension of a high-emissions scenario out to the year 2300, which is  
94 roughly eight times the pre-industrial CO<sub>2</sub> level and falls within the range that produced a good  
95 match to an Eocene-era climate in CESM1 (Zhu et al. 2019). Our approach combines a calculation  
96 of the high-latitude temperature distribution with a temperature budget along air parcel trajec-  
97 tories to analyze the influence of both initial conditions and diabatic evolution on cold air masses  
98 across the two scenarios. We observe a sharp drop-off in the wintertime temperature distribution  
99 at the freezing temperature, suppressing below-freezing conditions in the warmer climate and  
100 above-freezing conditions in the pre-industrial case. Using air parcel trajectories, we demonstrate  
101 a remarkable similarity in both the dynamics and diabatic effects acting on cold air masses across  
102 the two climate scenarios. Diabatic temperature evolution along cold air outbreak trajectories is a  
103 competition between cooling from longwave radiation and warming from boundary layer mixing  
104 of surface sensible heat flux. Surprisingly, while both diabatic effects strengthen in the warmer  
105 climate, the net effect remains the same, around  $-6$  K. By identifying the key physical processes  
106 influencing cold air outbreaks and how and why they will change in a warming climate, we hope  
107 to improve predictions of when and why North American cold air outbreaks can be expected to  
108 decline under anthropogenic climate change and demonstrate how Eocene-like climates could have  
109 maintained very warm continental interiors.

## 110 **2. Methods**

111 We approach the evolution of continental cold air outbreaks in a warming climate by comparing  
112 a pre-industrial to a warmer climate scenario using the Community Earth System Model, version 2

| Scenario       | Prescribed SST and Sea Ice  | Greenhouse Gases              |
|----------------|---|-------------------------------|
| Pre-industrial | From F1850  | CO <sub>2</sub> : 284.7 ppm   |
|                |   | CH <sub>4</sub> : 791.6 ppb   |
|                |   | N <sub>2</sub> O: 275.68 ppb  |
|                |   | F11: 12.48 ppt                |
| Year 2300      | From years 2290-2299 of the SSP5-8.5 fully coupled run of CESM2-WACCM from CMIP6, with minor corrections based on Hurrell et al. (2008) | F12: 0                        |
|                |   | CO <sub>2</sub> : 2166.15 ppm |
|                |   | CH <sub>4</sub> : 1070.33 ppb |
|                |   | N <sub>2</sub> O: 410.72 ppb  |
|                |   | F11: 333.48 ppt               |
|                |   | F12: 18.43 ppt                |

TABLE 1. Details of model setup for the pre-industrial and year 2300 warmer climate scenarios. Pre-industrial GHGs, SST, and sea ice are the same used by the F1850 component set. Year 2300 GHGs are the average over 2290-2300 of the extended SSP5-8.5 scenario (Meinshausen et al. 2020). Corrections to SST and sea ice (Hurrell et al. 2008) for Year 2300: SST =  $-1.8^{\circ}\text{C}$  if either SST  $< -1.8^{\circ}\text{C}$  or ice fraction  $\geq 90\%$ , ice fraction = 0 if  $< 0$  and = 100% if  $> 100\%$ , and ice fraction = 0 if SST  $> 4.97^{\circ}\text{C}$ .

(CESM2). Both scenarios start with a 2000-era climate with prescribed sea surface temperatures and sea ice (the F2000climo component set). For the pre-industrial scenario, we replace the prescribed year 2000 sea surface temperature and sea ice distributions with those of a pre-industrial case, circa 1850, and the atmospheric greenhouse gases with fixed 1850 concentrations. For the year 2300 warmer climate scenario, we prescribe the fixed greenhouse gas concentrations and the sea surface temperature and sea ice distributions using the average over the final decade from the fully coupled run of CESM2 under the SSP5-8.5 high-emissions scenario for CMIP6 extended out to the year 2300 (see Figures S1 and S2 in the supplementary material for surface temperature climatology of this model scenario). Details of both model scenarios can be found in Table 1. The first four years are discarded for spin-up, leaving 50 winters of data that are used in the analysis.

Our methodology for identifying cold air outbreaks, calculating trajectories, and interpolating meteorological variables onto those trajectories follows (Hartig et al. 2023). The reader is directed to that paper for additional detail on the methodology that follows.

We identify cold air outbreaks by sampling from the cold tail of a temperature distribution. To focus on continental cold air outbreaks, we define a sampling region over the interior of North America from 43 to 50°N and from 92 to 104°W, shown as a grey dashed box in Figures 2 and

134 4. Within the sampling region and for each model scenario, we generate a distribution of the  
135 wintertime (December through February) hourly 2-meter air temperature, shown in Figure 1. Cold  
136 air outbreaks are identified by randomly sampling from the coldest 5% of that distribution. To  
137 avoid double-counting a given air mass, we impose the additional constraint that samples must be  
138 at least three days apart from each other. Under this constraint, we sampled 400 cold air outbreaks  
139 in the pre-industrial scenario and 500 in the warmer climate.

140 We use air parcel backtracking to follow the evolution of air masses resulting in cold air outbreaks  
141 over North America. Each cold air outbreak identified with the method above is used to initialize  
142 a 10-day backwards trajectory, starting 100 meters above the ground at the location within the  
143 sampling region corresponding to the event. We perform the trajectory calculations with HYSPLIT,  
144 a trajectory and dispersion model developed by NOAA’s Air Resources Laboratory (Draxler and  
145 Hess 1998, 1997; Draxler 1999; Stein et al. 2015), with two custom modifications to improve  
146 data resolution described in detail in (Hartig et al. 2023). The trajectory calculations (resulting  
147 in latitude/longitude/altitude vs. time) are performed using hourly wind and other meteorological  
148 fields from the CESM2 model scenarios. Meteorological quantities are then interpolated onto the  
149 latitude, longitude and altitude of the air parcel for each hour along the trajectory.

150 Of particular interest in this study are the factors influencing temperature change within the  
151 model. The model temperature tendency at each gridpoint  $\dot{T}$  can be decomposed as,

$$152 \quad \dot{T} = \dot{T}_{physics} + \dot{T}_{dynamics} + \dot{T}_{fix}, \quad (1)$$

153 where  $\dot{T}_{physics}$  is the diabatic tendency from the model physics while  $\dot{T}_{dynamics}$  is calculated by  
154 the model’s dynamical core and accounts for advection, adiabatic compression and expansion, and  
155 the divergence damping (which is expected to be small). The term  $\dot{T}_{fix}$  is a very small energy  
156 correction, typically  $< 10^{-4}$  K/hr, that ensures the conservation of global energy.

157 Ultimately, we want to separate out the diabatic temperature tendencies  $\dot{T}_{physics}$  from the advective  
158 and adiabatic terms  $\dot{T}_{dynamics}$ . To remove the advective component of the model dynamics term,  
159 we use a Lagrangian reference frame by considering temperature tendencies along an air parcel  
160 trajectory instead of at a fixed point. To remove the adiabatic component, we transform from

161 temperature to dry static energy, which is conserved under adiabatic motion,

$$162 \quad DSE = c_p T + gz, \quad (2)$$

163 where  $T$  is temperature,  $g$  is gravitational acceleration,  $c_p$  is the specific heat of air, and  $z$  is  
164 geopotential height ( $z = 0$  at sea level). Whenever we refer to dry static energy throughout the rest  
165 of this paper, it is divided by the specific heat of air  $c_p$  to give units of temperature.

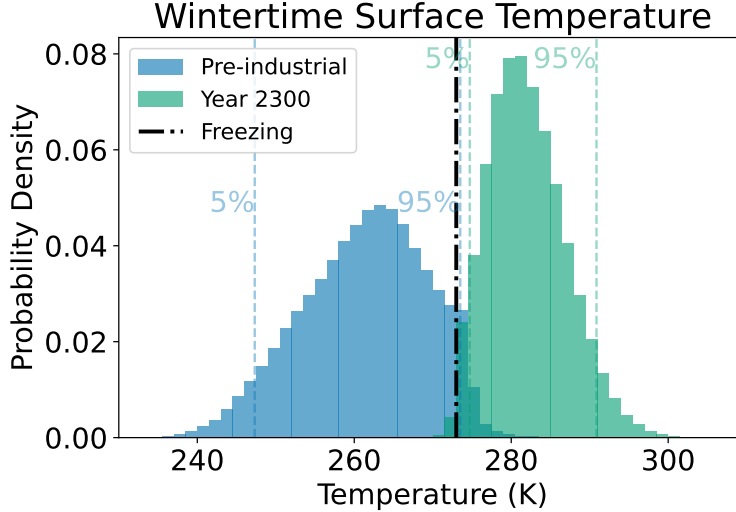
166 These transformations leave us with a dry static energy budget along air parcel trajectories  
167 consisting of temperature tendencies from distinct physics processes,

$$\begin{aligned} 168 \quad D\dot{S}E &= \dot{T}_{physics} \\ 169 \quad &= BL + DP + LW + SW + VD + GW + \dot{T}_{fix} \\ 170 \quad &\approx BL + DP + LW, \end{aligned} \quad (3)$$

172 where  $BL$  is boundary layer mixing,  $DP$  is deep convection,  $LW$  is longwave radiation,  $SW$  is  
173 shortwave radiation,  $VD$  is vertical diffusion, and  $GW$  is gravity wave drag. In the final line, we  
174 drop the last four terms, which are consistently over an order of magnitude smaller than the leading  
175 three terms in the troposphere over North America during winter. We also note that the tendencies  
176 due to boundary layer mixing  $BL$  and deep convection  $DP$  are shorthand for a partitioning of the  
177 moist processes temperature tendency within the atmosphere model,  $\dot{T}_{moist\ processes} = BL + DP$ , and  
178 are used as umbrella terms for a collection of processes.  $BL$  refers to the temperature tendency from  
179 the CLUBB parametrization (Golaz et al. 2002; Bogenschutz et al. 2013) and includes boundary  
180 layer turbulence, shallow convection, and latent heat from liquid cloud formation and evaporation,  
181 while  $DP$  is the remainder of moist processes and includes deep convection, cloud microphysics,  
182 and re-evaporation of rain and snow.

183 By integrating over time along a trajectory, we formulate a Lagrangian dry static energy budget  
184 that decomposes the total change in dry static energy of the air parcel into the diabatic contributions





196 FIG. 1. Narrowing of the wintertime surface temperature distribution in warmer climates. Distribution of  
 197 hourly 2-m surface air temperature over the interior of North America (see grey dashed box in Figure 2) between  
 198 December 1st and February 28th over 50 simulated winters. The Pre-industrial scenario is shown in blue and the  
 199 Year 2300 scenario in green. For reference, the 5th and 95th percentiles are marked by colored dashed lines and  
 200 the freezing temperature (273 K) by a black dashed line.

185 from distinct model physics processes,

$$\begin{aligned}
 \Delta DSE &\approx \int (BL + DP + LW) dt \\
 &= \int \dot{T}_{physics} dt.
 \end{aligned}
 \tag{4}$$

### 189 3. Results

190 We begin with overall changes to the wintertime temperature distribution over the interior of  
 191 North America between a pre-industrial and a much warmer climate in section a. The effect of  
 192 changes to the climatology and geographical distribution of source regions for cold air masses is  
 193 considered in section b, while the role of diabatic processes acting on those air masses as they are  
 194 swept over the interior of North America are covered in section c.

195 *a. Narrowing of the wintertime temperature distribution*

201 We begin with a consideration of the changes to the wintertime temperature distribution itself  
202 when moving from a pre-industrial climate to a much warmer climate. Figure 1 shows the Dec-Jan-  
203 Feb distributions of hourly 2-m air temperature over the sampling region in the interior of North  
204 America in the pre-industrial (blue) and year 2300 (green) modeled climate scenarios. In a warmer  
205 climate, the wintertime temperature distribution not only warms but also becomes narrower, with  
206 a cold tail that barely drops below freezing. The freezing temperature itself (vertical black line  
207 in the figure) helps to illustrate the dramatic change between the two scenarios: the wintertime  
208 temperature in this region almost never rises above freezing in the pre-industrial scenario and almost  
209 never drops below freezing in the warmer climate. There is, in fact, a particularly abrupt drop  
210 in both scenarios right around freezing. We argue that this suppression of temperature extremes  
211 around the freezing temperature is a result of the energy consumed by the latent heat of freezing in  
212 soil water. Using a range of volumetric heat capacities from  $1.5 \times 10^6$  to  $3 \times 10^6$  J/m<sup>3</sup>·K to account  
213 for variations in soil type and moisture content (Abu-Hamdeh 2003), we find with a back-of-the-  
214 envelope calculation (detailed in the supplementary materials and Figure S3) that the top 10 cm of  
215 soil contains more than enough water for the latent heat of freezing to entirely offset the surface  
216 energy imbalance that accumulates over the two days leading up to a typical cold air outbreak.  
217 Near-surface air temperatures would thus stall out at the freezing temperature as energy is diverted  
218 to melt or freeze water in the soil, reducing the prevalence of above-freezing temperatures in the  
219 pre-industrial climate and of below-freezing temperatures in the warmer climate. We note that this  
220 mechanism assumes efficient energy exchange between the soil surface and the lowest layers of  
221 air as well as within the soil to transmit the latent heat of freezing up into contact with overlying  
222 air, either of which might be exaggerated in the model to compensate for vertical discretization.  
223 Further work is required to confirm that the control that the phase change of soil water appears to  
224 have on near-surface air temperatures in CESM also operates in the real climate.

225 The increase in average wintertime temperature between the pre-industrial period and our year  
226 2300 scenario, as demonstrated by Figure 1, is unsurprising. Between a CO<sub>2</sub> concentration  
227 almost eight times that of the pre-industrial case and a total loss of Arctic sea ice, a wintertime  
228 mean temperature of 282 K (9°C, an increase of 21 K relative to pre-industrial) for the warmer  
229 climate is in line with paleoclimate proxies for the Eocene warm climate period (56–34 Myr),

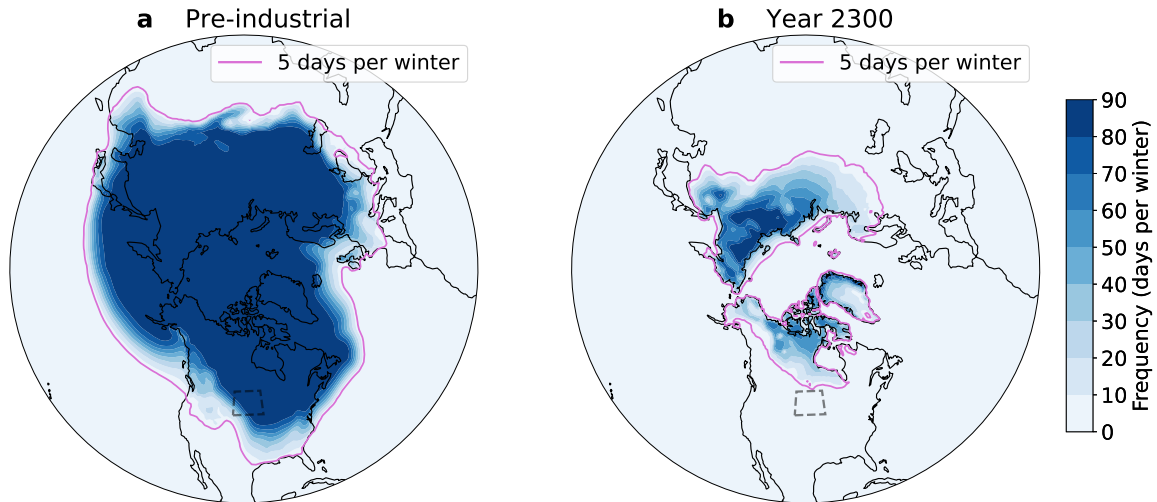
230 when frost-intolerant species such as crocodiles and palm fronds populated the interior of North  
231 America (Hutchison 1982; Wing and Greenwood 1993; Markwick 1994; Greenwood and Wing  
232 1995; Hyland et al. 2018). What is more interesting, at least from a dynamical perspective, is  
233 the change in temperature extremes. The 5th percentile of hourly temperature has increased by  
234 27 K between the two scenarios, or 1.3 times the increase in the wintertime mean, while the 95th  
235 percentile has increased by only 0.8 times the mean. The amplified warming of cold extremes is  
236 consistent with (Cronin and Tziperman 2015) who present a low cloud mechanism for suppressing  
237 cold air formation in a warmer climate. In our modeled scenarios, cold air outbreaks have indeed  
238 warmed by more than both the mean and hot extremes; understanding why this happens is a central  
239 goal of this paper.

240 To understand why wintertime temperature extremes are suppressed in a warmer climate, we  
241 consider changes in temperature in terms of two interrelated factors, (1) the location and temperature  
242 of source regions for cold air masses and (2) diabatic forcings acting on those air masses as  
243 they travel. The diabatic forcings include solar and infrared radiation, latent heating from the  
244 condensation or evaporation of water, and boundary fluxes like the surface sensible heat flux. In  
245 the shift to a warmer climate, changes to CO<sub>2</sub> concentration, atmospheric moisture, clouds, land  
246 cover, and precipitation can alter the diabatic forcings acting on an air mass as it travels. Source  
247 regions, on the other hand, can be thought of as setting an initial state of temperature and moisture  
248 that both is acted upon by and influences diabatic forcings. In our study of cold extremes over the  
249 interior of North America, we will therefore consider both changes to the climatology of source  
250 regions and changes to diabatic temperature tendencies along air mass trajectories in our attempt  
251 to tease out the causes of cold air suppression in a warmer climate.

### 252 *b. A warmer and less stable source region*

260 Beginning with temperature changes in source regions, we note the dramatic reduction in fre-  
261 quency of below-freezing days in Arctic source regions between the pre-industrial and the year  
262 2300 scenarios illustrated in Figure 2. The coldest temperatures over the interior of North America  
263 in the modern climate usually result from anomalous advection of air masses out of Arctic regions,  
264 where dryness and low temperatures set a cold initial state that is enhanced or maintained as these  
265 air masses cross the continent (Walsh et al. 2001; Cellitti et al. 2006; Portis et al. 2006; Vavrus et al.

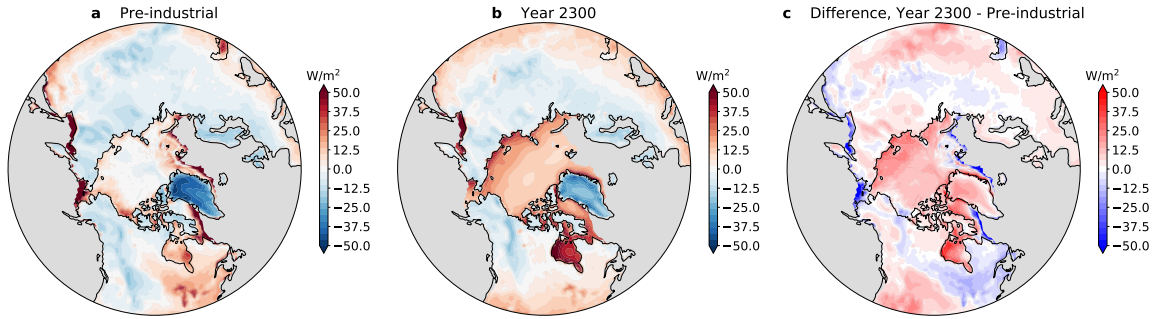
Days per winter when surface dry static energy is below freezing



253 FIG. 2. Reduction in area below freezing over continental cold air source regions in a warmer climate. Number  
254 of days per winter (December through February, averaged over 50 simulated winters) when the average daily  
255 surface dry static energy is above freezing. A solid contour has been included in fuchsia marking five days  
256 per winter to better demarcate the low-frequency regions. Note that, for this figure, dry static energy has been  
257 adjusted to a reference height of 510 m instead of sea level such that it reflects the temperature the air parcel  
258 would be if raised or lowered adiabatically to the average surface height over interior North America (grey dashed  
259 box in figure).

266 2006; Kolstad et al. 2010; Hanks and Walsh 2011; Smith and Sheridan 2018; Hartig et al. 2023).  
267 By using surface dry static energy relative to a reference height of 510 m instead of temperature, we  
268 are actually looking at the temperature these air masses would have after the adiabatic compression  
269 or expansion due to vertical motions that would accompany advection into the interior of North  
270 America (specifically, the grey dashed box in the figure, which has an average surface elevation  
271 of 510 m; for the same figure using surface air temperature, see Figure S4 in the supplementary  
272 materials). We leave the consideration of diabatic processes for the second half of this paper.  
273 Figure 2 shows a nearly order-of-magnitude reduction in the frequency of below-freezing days over  
274 the northernmost parts of North America and a total loss of such conditions over the newly exposed  
275 Arctic Ocean in the year 2300 scenario. We also note that Figure 2 is, if anything, a conservative  
276 estimate of the availability of below-freezing air masses for North America. Adiabatic compression  
277 will heat any air mass brought to the surface from aloft, so if air masses destined for North America

Winter climatology of surface sensible heat flux

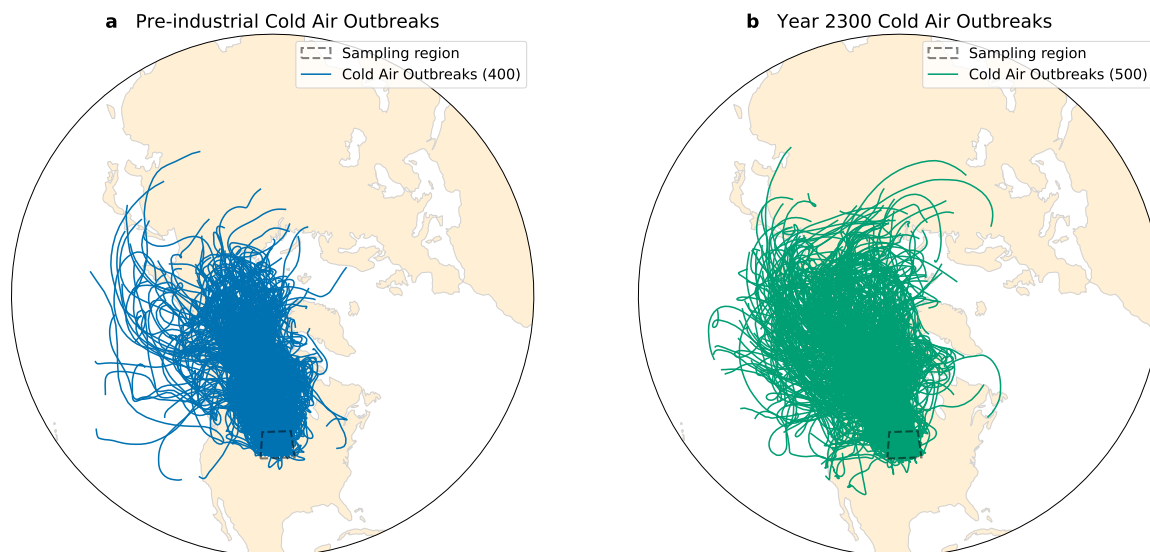


280 FIG. 3. Changes in surface sensible heat flux with the disappearance of Arctic sea ice. Wintertime climatology  
281 over 50 simulated winters of surface sensible heat flux, masked in grey to exclude regions with open ocean in the  
282 pre-industrial period (both sea ice and land fraction less than 50%). Includes Pre-industrial model climatology  
283 (left), Year 2300 (center), and the difference between the two scenarios (right). Positive values indicate heat flux  
284 from the surface into the atmosphere.

278 are sourced even a few hundred meters above the surface, then the heating from subsidence can  
279 partly or fully counteract the below-freezing initial temperature.

285 The biggest change in the year 2300 climate scenario other than the increase in CO<sub>2</sub> is the  
286 complete loss of Arctic sea ice and subsequent dramatic changes in heat and moisture fluxes in  
287 the Arctic. With the ocean surface fully exposed, surface sensible heat fluxes into the atmosphere  
288 increase by tens of W/m<sup>2</sup> relative to the pre-industrial climate, as shown in Figure 3. The loss  
289 of insulating sea ice makes near-surface air over the Arctic Ocean warmer rather than colder than  
290 bordering continental air, and provides a heat source to any air masses passing over it. For instance,  
291 while a large number of below-freezing days persist over Siberia in the warmer climate (Figure 2),  
292 any cold air masses originating there would need to pass over open ocean to reach North America,  
293 where they would be subject to significant positive surface heat fluxes (Figure 3).

300 Given the dramatic changes in the distribution of cold air in the northern high latitudes discussed  
301 above, we begin our air parcel trajectory analysis with an investigation into the extent to which  
302 the geographical distribution of source regions may have shifted in the warmer climate. After  
303 identifying cold air outbreaks over the interior of North America (grey dashed box), we calculate  
304 a 10-day back trajectory for each event (see section 2 for details). The resulting trajectory paths  
305 are shown in Figure 4 for both the pre-industrial and year 2300 model scenarios. While a larger

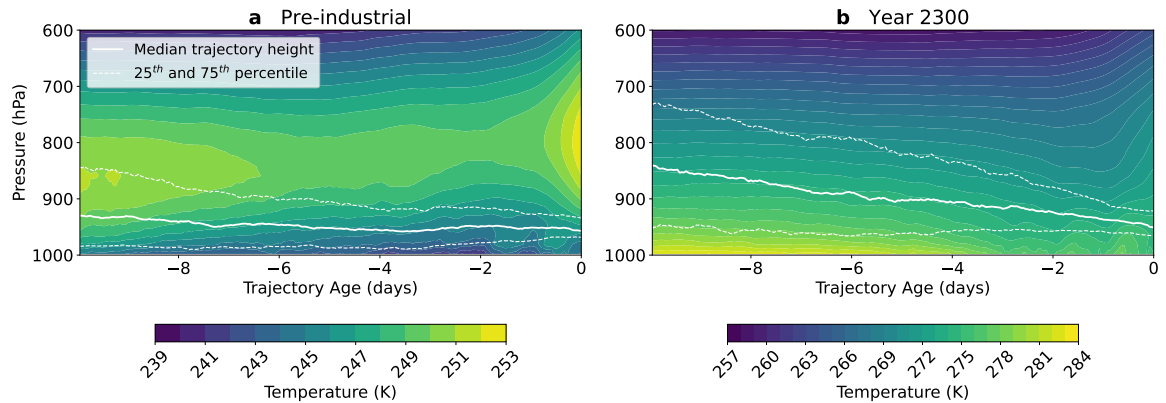


294 FIG. 4. Back trajectories of cold air outbreaks. Each line represents a 10-day back trajectory initialized from a  
 295 cold air outbreak identified in the pre-industrial (left, blue) or year 2300 (right, green) model scenario. The cold  
 296 air outbreaks used to initialize each trajectory were identified by randomly sampling 400 times (for pre-industrial  
 297 scenario) or 500 times (for year 2300 scenario) from the coldest 5% of hourly 2-m air temperatures over 50  
 298 simulated winters within the sampling region (black dashed box); see section 2 for more details on the sampling  
 299 method.

306 fraction of cold air outbreaks pass over the Pacific Ocean in the year 2300 case, it is remarkable how  
 307 little the source region and air mass trajectories have changed overall. Over half of all trajectories  
 308 still pass over the Arctic Ocean in the warmer climate scenario, even in spite of the positive heat  
 309 fluxes (Figure 3) and dramatically reduced availability of below-freezing days there (Figure 2).

317 Another climatological change in the northern high latitudes between the two climate scenarios  
 318 becomes obvious in the vertical temperature distribution, shown in Figure 5 as a composite of  
 319 air columns averaged over all cold air outbreak trajectories. In the pre-industrial climate, there is  
 320 a persistent near-surface temperature inversion in the winter high latitudes. The presence of the  
 321 inversion results in cold air outbreak trajectories that stay close to the ground; these near-surface  
 322 air masses are actually colder than the air aloft, and low or even negative sensible heat fluxes  
 323 along the trajectory paths (Figure 3) keep these air masses cold as they travel out of the Arctic.  
 324 In the warmer climate, the temperature inversion has disappeared in favor of a steady decline in  
 325 temperature with height. Positive surface sensible heat fluxes over the Arctic (Figure 3) in the

Composite of temperature profile along all trajectories



310 FIG. 5. Loss of the surface temperature inversion along cold air outbreak trajectories. Composites of the  
 311 vertical temperature profile over all cold air outbreak trajectories (colored lines in Figure 4) in the pre-industrial  
 312 (left) and year 2300 (right) model scenarios. The x-axis is the time along the trajectories in days, where day 0  
 313 represents the occurrence of the cold air outbreak in the sampling region. The spread of air parcel trajectory  
 314 heights within these composite profiles is indicated by a solid white line for the median height and dashed white  
 315 lines for the 25<sup>th</sup> – 75<sup>th</sup> percentile range across all trajectories. Note that the contour range is different between  
 316 the two figures, but the contour spacing is the same (1 K).

326 first half of the trajectories and cooler temperatures aloft than at the surface mean that cold air  
 327 outbreak air masses originate higher up in the atmosphere in the warmer climate than they do in  
 328 the pre-industrial case.

329 *c. Enhanced diabatic heating and cooling*

330 With the initial temperature determined by the climatology in source regions, diabatic heating  
 331 and cooling provide the final piece to determine temperature evolution as those air masses travel  
 332 into the mid-latitudes. If there is net diabatic heating along a trajectory, it can offset the initially  
 333 low temperatures of the Arctic air parcel and suppress cold air formation. Conversely, if there is  
 334 net diabatic cooling, the air mass can become even colder as it travels into the mid-latitudes, so  
 335 diabatic sources are critically important in determining the development of cold air outbreaks.

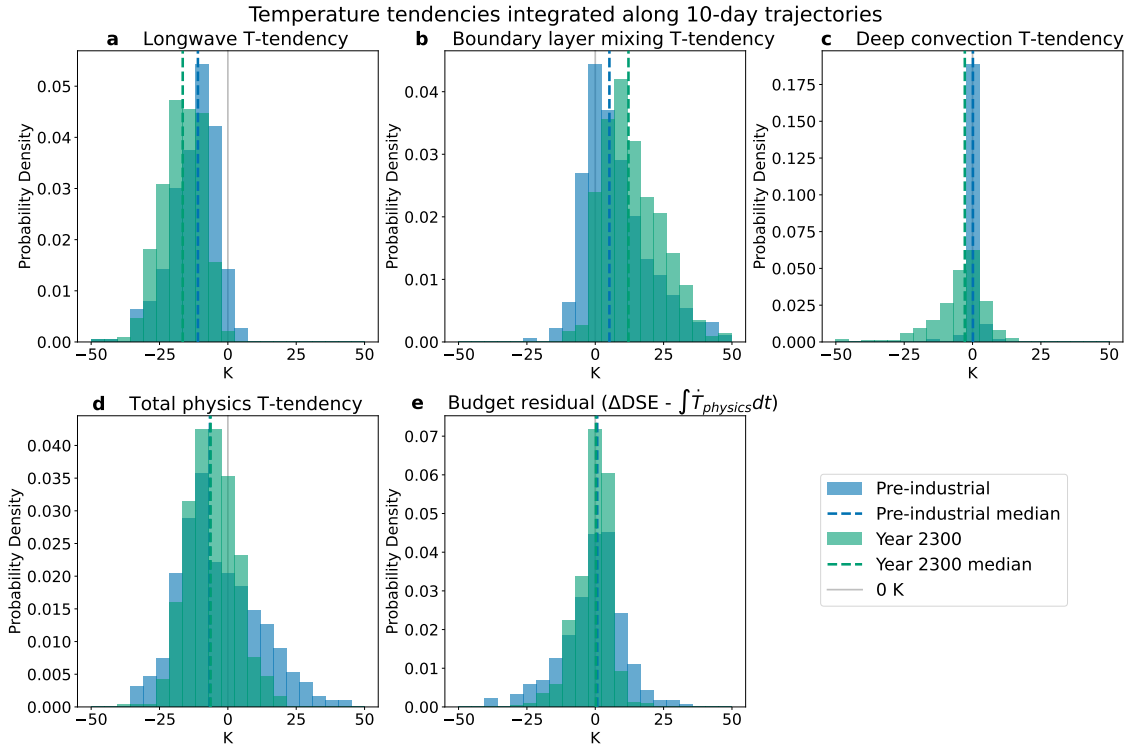
336 Using a Lagrangian air parcel perspective, we can separate the temperature tendency due to  
 337 diabatic forcings from the advective and adiabatic components. We take advantage of the way  
 338 that CAM6 separates temperature tendency into distinct physical processes to decompose diabatic

339 temperature evolution into five terms: boundary layer mixing, deep convection, longwave radiation,  
340 shortwave radiation, and gravity wave drag. The last two terms, shortwave radiation and gravity  
341 wave drag, are more than an order of magnitude smaller than the other components, so we leave them  
342 out of the analysis that follows (see Figure S7 in the supplementary materials for all components).  
343 Additionally, the tendencies due to boundary layer mixing and deep convection are each shorthand  
344 for a collection of processes within the atmosphere model that are explained in more detail in  
345 section 2. We interpolate each temperature tendency onto the air parcel trajectories and then  
346 integrate along each trajectory to get a distribution over all trajectories of the ten-day diabatic  
347 temperature tendency attributable to each physical process, shown in Figure 6.

356 We note a crucial balance in the diabatic temperature evolution of cold air masses in both the  
357 pre-industrial and year 2300 climate scenarios. In both scenarios, diabatic temperature evolution is  
358 almost entirely dominated by a competition between cooling from longwave radiation (Figure 6a)  
359 and warming from boundary layer mixing (Figure 6b). The key difference between the two  
360 scenarios is that both processes become more intense in the warmer climate; longwave cooling  
361 of the air mass is actually stronger in the warmer climate, while boundary layer mixing provides  
362 more heating at the parcel level. The two processes nearly cancel out when averaged over all  
363 trajectories, but the spread across trajectories is large (see Figure 6d), indicating that individual  
364 trajectories may experience large diabatic temperature change in either the positive or negative  
365 direction. For reference, we note that the 5th percentile of surface temperature is only 7.3 K below  
366 the winter average temperature of 282 K over the interior of North America in the year 2300  
367 scenario (Figure 1), and the surface temperature in source regions is only below freezing a few days  
368 every year (Figure 2). And while the *average* diabatic temperature change is just a few degrees, the  
369 large spread in Figure 6 means that most trajectories in fact have a non-zero diabatic temperature  
370 change. This means that the diabatic temperature change, which is commonly 10 or 20 K over ten  
371 days according to Figure 6, is sufficient to entirely determine whether an Arctic air mass becomes  
372 a cold air outbreak in the year 2300 scenario, and to a lesser extent in the pre-industrial scenario as  
373 well.

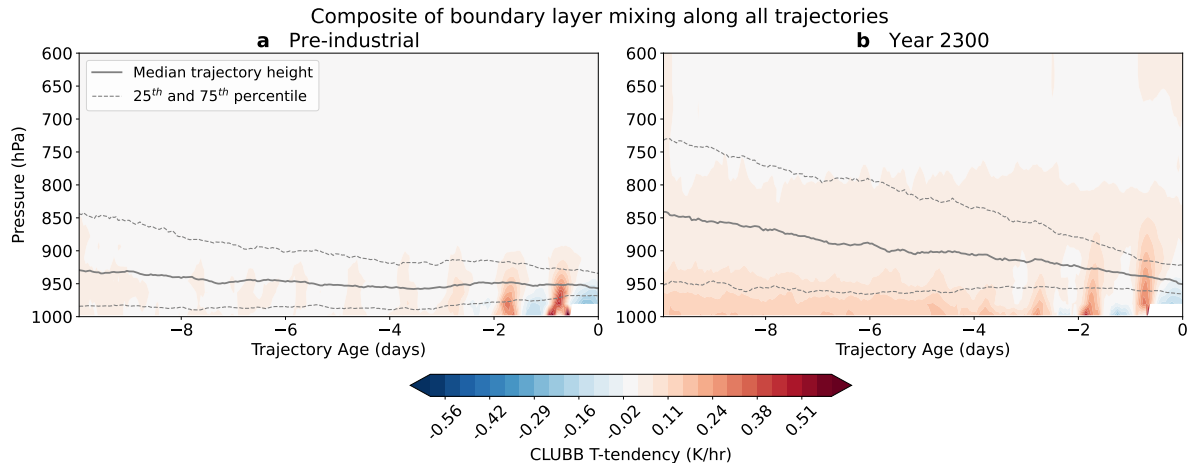
374 Two questions naturally arise from the changes in diabatic temperature contributions between  
375 the pre-industrial and year 2300 scenarios in Figure 6a and 6b: why does boundary layer mixing  
376 lead to more heating in the warmer climate? And why is longwave cooling more intense?





348 FIG. 6. Contributions to diabatic temperature evolution along cold air outbreak trajectories. Each histogram  
 349 shows the distribution across all trajectories of the cumulative temperature contribution of a specific model  
 350 physics process, found by integrating the hourly temperature tendency along each 10-day trajectory. The three  
 351 largest contributions to the total diabatic temperature evolution are shown along the top row: longwave radiation  
 352 (a), boundary layer mixing and latent heat of condensation/evaporation (b), and deep convection and cloud  
 353 microphysics (c). The bottom row shows the total diabatic temperature tendency (d) and the total budget closure  
 354 residual (e;  $\Delta DSE - \int \dot{T}_{physics} dt$ ), which measures how well changes in dry static energy match up with the  
 355 diabatic temperature tendencies given by the model.

381 There are two key changes in the warmer climate that can facilitate both heating from boundary  
 382 layer mixing and cooling from deep convection. The loss of the surface temperature inversion  
 383 (Figure 5) removes a major barrier to convection, allowing both shallow and deep convection to  
 384 proceed more readily. The increase in surface fluxes of heat (Figure 3) and moisture into the  
 385 atmosphere inject more heat into the boundary layer and can help fuel convection. Mixing in  
 386 the boundary layer redistributes those surface fluxes throughout the boundary layer, leading to  
 387 heating as shown in Figure 7, while deep convection transports surface heat upwards into the  
 388 mid-troposphere, leading to cooling in the boundary layer and heating aloft (see Figure S8 in the

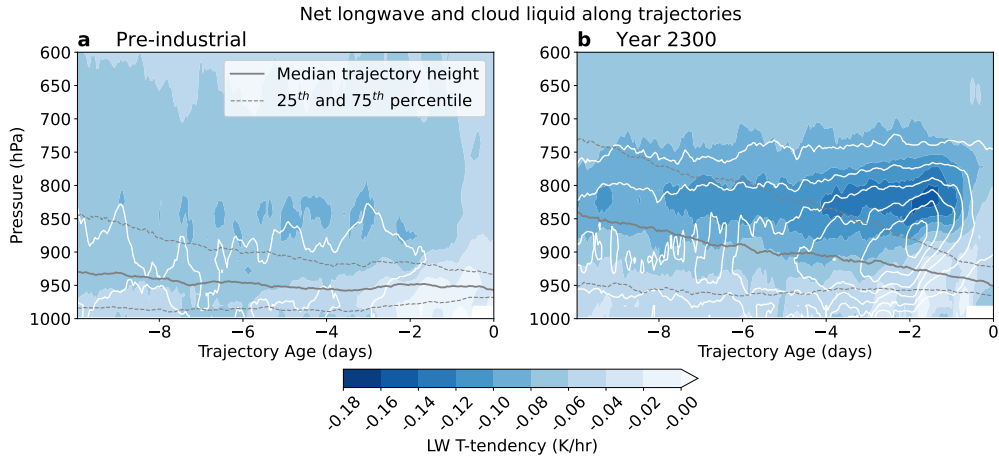


377 FIG. 7. Boundary layer mixing brings more heat higher up in the warmer climate. Composites of the boundary  
 378 layer mixing temperature tendency profiles (shaded contours) over all cold air outbreak trajectories (colored  
 379 paths in Figure 4) in the pre-industrial (left) and year 2300 (right) model scenarios. Grey lines mark the median  
 380 (solid) and 25<sup>th</sup>–75<sup>th</sup> percentile range (dashed) of trajectory heights across all trajectories.

389 supplementary materials for deep convection profiles). Evidently, more heat is distributed within  
 390 the boundary layer where most of these trajectories reside than is carried aloft (cooling from deep  
 391 convection in Figure 6c for most trajectories is weaker than heating from boundary layer mixing  
 392 in 6c), but both processes are enhanced by the loss of the near-surface stable layer and increased  
 393 surface heat fluxes in the migration from a pre-industrial to a warmer climate.

399 Longwave radiation is the only significant source of diabatic cooling that remains in the warmer  
 400 climate. Deep convection is just  $-5$  K over ten days on average in the warmer climate, while  
 401 boundary layer mixing becomes almost strictly positive. Continental cold air outbreaks, therefore,  
 402 rely almost entirely on longwave cooling to make or keep Arctic air masses cold into the mid-  
 403 latitudes. The intensification of longwave cooling in the warmer climate is therefore crucial to  
 404 identifying how cold air outbreaks respond to warming.

405 We find that the increase in longwave cooling is due primarily to cloud-top radiative cooling  
 406 between 800 and 900 hPa, where cold air outbreaks in the warmer climate spend much of their  
 407 time. Figure 8b demonstrates this effect with composites of cloud liquid and the vertical structure  
 408 of the longwave temperature tendency along trajectories. In the warmer climate, cloud liquid  
 409 peaks around 900 hPa and the longwave temperature tendency peaks along the top of the cloud  
 410 layer between 800 and 850 hPa at values around 0.1 K/hr, or upwards of 2 K/day. In the pre-



394 FIG. 8. Cloud-top radiative cooling enhances longwave diabatic cooling in the warmer climate scenario.  
 395 Composites of the longwave temperature tendency (shaded contours) and cloud liquid (solid contours) profiles  
 396 over all cold air outbreak trajectories in the pre-industrial (left) and year 2300 (right) model scenarios. Grey lines  
 397 mark the median (solid) and 25<sup>th</sup>–75<sup>th</sup> percentile range (dashed) of trajectory heights across all trajectories.  
 398 Cloud liquid contours start at  $1 \times 10^{-5}$  kg/kg and have a constant spacing of  $1 \times 10^{-5}$  kg/kg.

411 industrial climate, by contrast (Figure 8a), cloud liquid is almost an order of magnitude lower  
 412 and the longwave temperature tendency does not vary much with height. The presence of thicker  
 413 liquid clouds, made possible by higher temperatures and moisture in the warmer climate, increases  
 414 longwave radiative cooling within the boundary layer and is responsible for more intense longwave  
 415 radiative cooling along cold air outbreak trajectories.

#### 416 4. Discussion and Conclusions

417 Wintertime cold air outbreaks affect large swathes of the interior of North America, but their  
 418 behavior under global warming has proven challenging to predict. In spite of the mean warming  
 419 trend over the last few decades and its amplification in the Arctic, there is disagreement on  
 420 whether cold air outbreaks have declined as well. We know that warming wins out eventually,  
 421 because warmer paleoclimate periods like the Eocene (56–34 Mya) present fossil evidence of  
 422 strong suppression of cold extremes over continental interiors. Therefore, a better understanding  
 423 of the mechanisms that sustain or suppress cold extremes in a variety of climate states could bolster  
 424 our knowledge of paleoclimates and improve predictions of temperature extremes under future  
 425 warming.

426 In this study, we analyze the development of North American wintertime cold air outbreaks using  
427 the Community Earth System Model, CESM2. We compare two climate scenarios with prescribed  
428 greenhouse gases, sea surface temperature, and sea ice coverage, corresponding to a pre-industrial  
429 case and a high-emissions case circa the year 2300, which produces Eocene-like conditions. A  
430 mid-latitude cold air outbreak may involve some combination of cold initial temperatures in the  
431 source region and cooling along the path of travel due to diabatic effects, so we consider both of  
432 these factors in our analysis. In the first half of the paper, we analyze climatological differences in  
433 cold air outbreak source regions between the pre-industrial and the much warmer climate scenario  
434 to quantify the availability of cold air. In the second half, we use a breakdown of the temperature  
435 tendencies due to distinct physics processes affecting air parcels as they travel into the mid-latitudes  
436 to clarify the role of diabatic heating and cooling in turning Arctic air masses into mid-latitude  
437 cold air outbreaks.

438 We began this paper by posing two key questions: given the persistence of North American cold  
439 air outbreaks observed over the last few decades in spite of an overall warming trend, when and why  
440 will they decline as warming continues? And how did much warmer past climates like the Eocene,  
441 which may serve as an analog for warmer future climates, so effectively suppress below-freezing  
442 temperatures over continental interiors? In the process of addressing these questions below, we  
443 highlight our three main results: (1) the latent heat of freezing suppresses warm extremes in a  
444 pre-industrial climate and cold extremes in a much warmer climate; (2) the transition to a warmer  
445 climate, and in particular the loss of Arctic sea ice and the near-surface temperature inversion  
446 at high latitudes, dramatically decreases the availability of below-freezing air, suppressing mid-  
447 latitude cold extremes; and (3) while the net diabatic temperature change along cold air outbreak  
448 trajectories is nearly identical between the pre-industrial and the much warmer climate, the primary  
449 heating and cooling mechanisms both get stronger in the warmer climate.

450 Based on wintertime temperature distributions over the interior of North America, we determined  
451 that the latent heat of freezing serves as a significant barrier to extreme temperatures in both climate  
452 scenarios. In a warmer climate, where wintertime surface temperatures are generally above  
453 freezing, liquid water in surface soils has the potential to release enough heat through freezing  
454 to offset the energy lost to surface sensible heat flux as a cold air mass approaches, preventing  
455 surface temperatures from dropping below freezing. The converse occurs in a cold climate, as

456 ice near the surface can absorb excess heat through melting and resist the formation of above-  
457 freezing air. Both effects are visible in the sharp drop of the wintertime temperature distribution  
458 at freezing (Figure 1), shortening the warm tail of the pre-industrial distribution and suppressing  
459 below-freezing temperatures in the warmer climate. This result implies that the availability of  
460 surface and soil water can have an important mitigating effect on the development of temperature  
461 extremes near the freezing temperature.

462 The transition to a warmer climate is also accompanied by a marked decrease in the availability of  
463 below-freezing air in the Arctic source region, a result of changes to a handful of key climatological  
464 features. Increased heat fluxes from the ocean to the atmosphere at high latitudes (Figure 3), a  
465 result of the disappearance of Arctic sea ice in the warmer climate, raise surface temperatures  
466 over the Arctic Ocean above freezing, dramatically reducing the availability of below-freezing air  
467 near the surface (Figure 2). The near-surface temperature inversion, present throughout the winter  
468 at high latitudes in the pre-industrial climate, also disappears in the warmer climate (Figure 5),  
469 a manifestation of Arctic amplification. This removes the layer of stable, cold, dry air near the  
470 surface that supplied most of the cold air masses in the pre-industrial case. It also presents an  
471 avenue to predicting the decline of cold air outbreaks in the near future: Arctic sea ice and the  
472 near-surface temperature inversion are crucial to supplying the extremely cold air masses that can  
473 turn into mid-latitude cold air outbreaks. Disruptions to either of those features are likely to result  
474 in a significant shift in the availability of cold air masses.

475 Given the dramatic changes to temperature and stability in source regions that accompany the  
476 transition to a warmer climate, it is remarkable how little the dynamics and net diabatic effects  
477 were found to change. Cold air outbreak trajectories are less heavily clustered over the Arctic  
478 Ocean in the warmer climate relative to the pre-industrial case, but over half still pass over on their  
479 way to the mid-latitudes (Figure 4). The Arctic therefore continues to serve as the primary source  
480 region for cold air outbreaks. The net diabatic temperature change along trajectories is also nearly  
481 identical, at  $-6.5$  K for the pre-industrial and  $-6.3$  K for the warmer climate.

482 The similarity in net diabatic effect across the two climate scenarios would be surprising in its  
483 own right considering all of the changes to atmospheric temperature and moisture that accompany  
484 the shift to an Eocene-like climate, but is even more remarkable given the changes to two of the  
485 largest contributors to diabatic temperature change, longwave radiation and boundary layer mixing

486 (Figure 6). Longwave cooling is more intense in the warmer climate due to enhanced cloud-top  
487 radiative cooling (Figure 8), while warming from boundary layer mixing is also stronger and re-  
488 flects more effective redistribution of enhanced surface sensible heat fluxes (Figure 7). While the  
489 net effect is the same for the pre-industrial and the warmer climate, the distribution of total physics  
490 temperature tendency is narrower (Figure 6d), which could reflect a greater correlation between  
491 the two competing effects via convection. This result also has implications for the challenges  
492 encountered by climate models attempting to represent cold air suppression in Eocene-like conti-  
493 nental interiors. The enhancement of both the longwave and boundary layer mixing temperature  
494 tendencies in the warmer climate is primarily a result of cloud and convection processes, which  
495 are notoriously tricky to simulate. Models may be missing or misrepresenting key elements of  
496 these processes in the extreme conditions of a much warmer climate, where we have few modern  
497 analogs with which to develop and test parameterizations.

498 When considering the scope and implications of this work, it is important to keep in mind several  
499 caveats. We note that our warmer climate scenario, intended to emulate the year 2300 of an  
500 extended high-emissions run of CESM, is not a direct simulation of the Eocene itself but rather  
501 evocative of much warmer climates. By using a 2000-era model configuration as a starting point,  
502 we incorporate modern vegetation, land ice, and topography while setting greenhouse gases to  
503 match projected anthropogenic emissions. We chose these conditions in pursuit of a signal in the  
504 occurrence and characteristics of cold events that is strong enough to be differentiated from that  
505 of the pre-industrial scenario and to allow for attribution to changes in radiative forcing, sea ice,  
506 and sea surface temperature rather than variable vegetation and topographic effects, which we feel  
507 was largely successful. There are also limitations to the interpretability of our trajectory analysis,  
508 which are also discussed in detail in Hartig et al. (2023) but will be restated here. The diabatic  
509 temperature budget residual for a Lagrangian trajectory, which should be zero, is near zero for the  
510 majority of trajectories but upwards of 10 K for a handful in both climate scenarios (Figure 6e).  
511 We attribute these residuals (budget errors) primarily to the effect of vertical positioning errors in  
512 HYSPLIT that are further exaggerated by the strong vertical temperature gradients in near-surface  
513 winter-time inversions. This is consistent with the reduced error in the warmer climate scenario  
514 where vertical temperature gradients are smaller due to the loss of the near-surface temperature  
515 inversion. As the median residual for both scenarios is zero, our focus remains on the differences

516 in temperature tendency between the two scenarios and the physical mechanisms we can identify  
517 that explain those differences. With our method, we are able to quantify the residual itself rather  
518 than attributing it to a collection of physical processes that do not have an explicit temperature  
519 tendency as is more commonly done.

520 Ultimately, we have developed a mechanistic understanding of North American cold air outbreaks  
521 and how they might change in a warmer climate. Both the availability of cold air in high-latitude  
522 source regions and diabatic effects acting to heat or cool air masses as they travel into the mid-  
523 latitudes can affect the development of wintertime cold extremes. The median diabatic cooling  
524 is  $-6$  K for air parcels traveling toward cold events. But it ranges into tens of degrees for many  
525 trajectories, so diabatic effects often play a major role in the evolution of Arctic air masses in both  
526 the pre-industrial and the much warmer climate scenario. Since the net diabatic effect is the same  
527 in the two scenarios, we find that the decreased availability of cold air in the Arctic in a warmer  
528 climate is sufficient to suppress below-freezing temperatures over the interior of North America  
529 in spite of the persistence of diabatic cooling along air parcel trajectories in the warmer climate  
530 scenario. To understand and predict changes to cold air outbreaks in a warming world, we must  
531 therefore account for changes to both source regions and diabatic processes, as either can tip the  
532 scale in controlling whether an Arctic air mass becomes a cold air outbreak when swept into the  
533 mid-latitudes.

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542 *Data availability statement.* Simulations were performed using CESM2.1.3 with CAM6, which  
543 is freely available at <https://www.cesm.ucar.edu/models/cesm2> following registration. The  
544 CESM2-WACCM output used to generate prescribed SST and sea ice files for the Year 2300  
545 scenario is available at <http://doi.org/10.22033/ESGF/CMIP6.10115> under the ssp585 tag.  
546 Fixed greenhouse gas concentrations for the Year 2300 scenario come from Meinshausen et al.  
547 (2020). Back trajectories were calculated using HYSPLIT v5.1.0 (with modifications described in  
548 Hartig et al. (2023)), which is freely available at <https://www.ready.noaa.gov/HYSPLIT.php>.  
549 Code used to interpolate CAM data onto HYSPLIT trajectories is available on GitHub at <https://github.com/kahartig/camtrack>.  
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